

Statistical-mechanical forcing of ocean circulation

Bill Merryfield

Canadian Centre for Climate Modelling and Analysis

Origins

Salmon, R., G. Holloway and M. C. Hendershott, 1976: The equilibrium statistical mechanics of simple quasi-geostrophic models. *J. Fluid Mech.*, 75, 691-703.

$$\frac{\partial q}{\partial t} + \mathbf{u} \cdot \nabla q = 0, \quad q = \nabla \times \mathbf{u} + h$$

$$h = \frac{f(H_o - H)}{H_o}$$

Conserves *E*, moments of $Q_n = \int q^n dA$ (circulation, enstrophy...)

Origins

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- \bullet Inviscid \rightarrow cons of energy, enstrophy
- Spectrally truncated
- Equilibrium via maximization of entropy

"...equilibrium...flow is positively correlated with bottom topography (anticyclonic flow over seamounts)"

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Example: circulation over bumpy ridge

- Consider 100 realizations from random initial conditions
- Ensemble mean flow develops on \sim eddy turnover time



pseudowestward (shallower water on right in NH)

Physics is advective rearrangement of PV, constrained by conservation of E, moments of q



What is Entropy?

 Entropy is the INFORMATION DEFICIT between detailed knowledge (microstate) and statistical knowledge (macrostate) of a system

Simple example

- *N* particles in a partitioned box: $N = n_L + n_R$
- $N=2 \rightarrow 2^2=4$ possible configurations



Simple example

- N particles in a partitioned box: $N = n_L + n_R$
- $N=100 \rightarrow 2^{100}$ possible configurations



Simple example

• Now introduce *time dependence*:

probability X (unit time)⁻¹ a particle will undergo *transition* between two boxes



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- dS/dt ≥ 0 in ensemble mean → arrow of time
- Maximizing S subject to constraints ($n_L + n_R = N$ here) locates equilibrium

To define *microstates*, represent *p* by sampling of *N* realizations (*N* sufficiently large to resolve finest scales of *p*)



• To define *macrostates*, divide phase space into *M* observably distinguishable cells, consider *observable* PDF $P_i = n_i / N$



Multiplicity (# microstates for given macrostate) is

 $w = \frac{N!}{\prod_i n_i!}$

 Corresponding *information deficit* or *entropy* is

 $S = k \log w$

(*k* = 1/log 2 if in bits)



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• Using $\log n_i \approx n_i \log n_i - n_i$, $P_i = n_i / N$

 $S = -k \sum_{i} P_{i} \log P_{i}$

or

$$S = -k \int P \log P \, dY$$



Salmon, R., G. Holloway and M. C. Hendershott, 1976: The equilibrium statistical mechanics of simple quasi-geostrophic models. *J. Fluid Mech.*, 75, 691-703.

Bretherton, F. P. and D. Haidvogel, 1976: Two-dimensional turbulence above topography. *J. Fluid Mech.*, 78, 129-154.

- Inviscid \rightarrow cons of energy, enstrophy
- Spectrally truncated
- Equilibrium via maximization of entropy

- Small-scale dissipation dissipates enstrophy, conserves energy
- No explicit truncation
- Equilibrium via minimization of enstrophy

"...equilibrium...flow is positively correlated with bottom topography (anticyclonic flow over seamounts)"

"...initially turbulent flow tends to a steady state with streamlines parallel to contours of constant depth, anticyclonic around a bump"

$$\mu\psi = q$$

Salmon, R., G. Holloway and M. C. Hendershott, 1976: The equilibrium statistical mechanics of simple quasigeostrophic models. *J. Fluid Mech.*, 75, 691-703.

Bretherton, F. P. and D. Haidvogel, 1976: Two-dimensional turbulence above topography. *J. Fluid Mech.*, 78, 129-154.

Carnevale, G. F. and J. S. Frederiksen, 1987: Nonlinear stability and statistical mechanics of flow over topography. *J. Fluid Mech.*, 175, 157-181.

"...in the limit of infinite resolution the canonical mean state is statistically sharp, that is, without eddy energy on any scale, and is identical to the nonlinearly stable minimum enstrophy state"

SHH76

"We suggest that some of the statistical trends observed in nonequilibrium flows may be looked on as manifestations of the tendency for turbulent interactions to maximize the entropy of the system."

World Ocean









World Ocean



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Equilibrium→ Disequilibrium





World Ocean



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Equilibrium→ Disequilibrium

Holloway (*JFM* 1978) Carnevale, Frisch & Salmon (*J Phys* 1981) Moment closure implies dS/dt≥0



Treguier (GAFD 1989)
Merryfield & Holloway (JFM 1997) $\mu < \psi > = < q >$ continues
to ~hold if eddy timescales <
forcing, dissipative timescales

Simple example revisited

- Remove constant-N constraint by adding A particles s⁻¹ to the left compartment, and extracting particles from the right compartment at the same rate
- This external forcing prevents equilibrium from being attained. However, the *tendency* for S increase due to random "dynamics" persists
- entropy increase tendency acts as "force" which balances applied force in statistical equilibrium





World Ocean









World Ocean









SHH76

World Ocean



World Ocean





























x (wm)

400

200

Barotropic streamfunction over topography

0

204

200

Merryfield, Cummins and Holloway (JPO 2001)

0 5 10 Streamfunction (10⁷ m³ s⁻¹)

<Ψ>

-5

15

Equilibrium flow along a shelf



Merryfield, Cummins and Holloway (JPO 2001)

Equilibrium flow along a shelf

→ Maximum current *upslope* of maximum slope

Unstratified



Merryfield, Cummins and Holloway (JPO 2001)

Equilibrium flow along a shelf

→ Maximum current *upslope* of maximum slope

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With depth dependence as in QG



Merryfield, Cummins and Holloway (JPO 2001)

Rectified flow over a seamount



Beckmann and Haidvogel (JGR 1997)

SHH76

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Observations







Example of eddy-active ocean model: Ocean Model for Earth Simulator (OFES)



Sea Surface Temperature (Day: 258, Year: 50)



∆**x=1/10**°

Topostrophy in OFES 1/10° model



Merryfield & Scott (Ocn Mod 2007)

Topostrophy in eddying vs non-eddying global ocean models





Observed topostrophy from current meter archive

Topostrophy vs. latitude and relative depth



Holloway (JGR 2008)

Observed topostrophy vs models sampled at current meter locations

Model	Number of points	Observed topostrophy	Modeled topostrophy
CCSM3	6921	0.302	0.116
CGCM3	5663	0.292	-0.026
OFES	6526	0.261	0.139
POP	6975	0.314	0.147

Eddy active models

Holloway (JGR 2008)

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How much resolution is "enough"?

Holloway (JGR 2008)



Model configuration



u over seamount



 \rightarrow .04 ms⁻¹ (all)

Whole-domain topostrophy













Ratio of abyssal to upper ocean KE KE_{>1000m} / KE_{<1000m}



What about parameterization?



• Back to arrow of time

 Holloway (JPO 1992): replace viscous operator with forcing toward higher entropy state

 Consequences of entropy-gradient forcing are not always obvious!









Intrepretation

- Consider ensemble mean
- Heuristically replace nonlinear terms by tendency toward equilibrium q^{*}_{i,k}:

$$\begin{split} &\frac{\partial \overline{q}_{1,k}}{\partial t} = -A_k (\overline{q}_{1,k} - \overline{q}_{1,k}^*), \\ &\frac{\partial \overline{q}_{2,k}}{\partial t} = -A_k (\overline{q}_{2,k} - \overline{q}_{2,k}^*) + v_0 k^2 \overline{\psi}_{2,k}. \end{split}$$

- Substitute equilibrium $q_{i,k}^*$
- Find ψ_i as a function of k, v_o/A

Increasing bottom friction

Increasing bottom friction

Rectified flow over a seamount

Beckmann and Haidvogel (JGR 1997)

Summary

- Equilibrium flow over topography generally in pseudowestward direction
- Non-ideal fluids experience "force" toward higher entropy
- Numerical and observational evidence is mounting that this effect has strong influence on (deep) ocean circulation
- OGCMs not converged, even at 1/10°

Statistical dynamics in the cosmos

Nonlinear Dynamics and Statistical Theories for Basic Geophysical Flows

Anderey J. Malda and Xinsmung Wang.

2d turbulence and stellar systems obey similar statistical mechanics

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2d turbulence and stellar systems obey similar statistical mechanics

 Π

After P.H. Chavanis

Picturing entropy

$$S = <\log P > = \frac{1}{2}\sum_{\mathbf{k}}\log E(\mathbf{k})$$

