
Stochastic Parameterisation Schemes Based on Rigorous Limit Theorems

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- Most stochastic climate models are specific to the modeled system, but system-specific SDE models in particular implicitly apply limit theorems, which give general formulas

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→ the equations of motion are linearised about a mean state,

$$\frac{dX}{dt} = LX,$$

white noise is added to account for the (fast-evolving) error in linearisation, and a damping term is added for stability:

$$\frac{dX}{dt} = (L + D)X + \frac{dW}{dt}$$

Theorem-based reduction methods

Hasselmann(1976)

Kurtz(1973) \longrightarrow MTV(1999-2006) \longrightarrow Franzke et al.(2005)

Papanicolaou(1976)

Khasminskii(1966)

Hasselmann(1976) \longrightarrow Arnold et al.(2003)

Fatkullin and Vanden-Eijnden(2004)

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Hasselmann(1976) \longrightarrow . . .

$$\begin{aligned}\frac{dx}{dt} &= f(x, y) && \text{(slow climate mode)} \\ \frac{dy}{dt} &= \frac{1}{\epsilon}g(x, y) && \text{(fast weather mode)}\end{aligned}$$

As $\epsilon \rightarrow 0$, $x \rightarrow X$ in distribution, where X satisfies:

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★ **Simple to implement: do not have to resolve fast mode, y**

Online closure

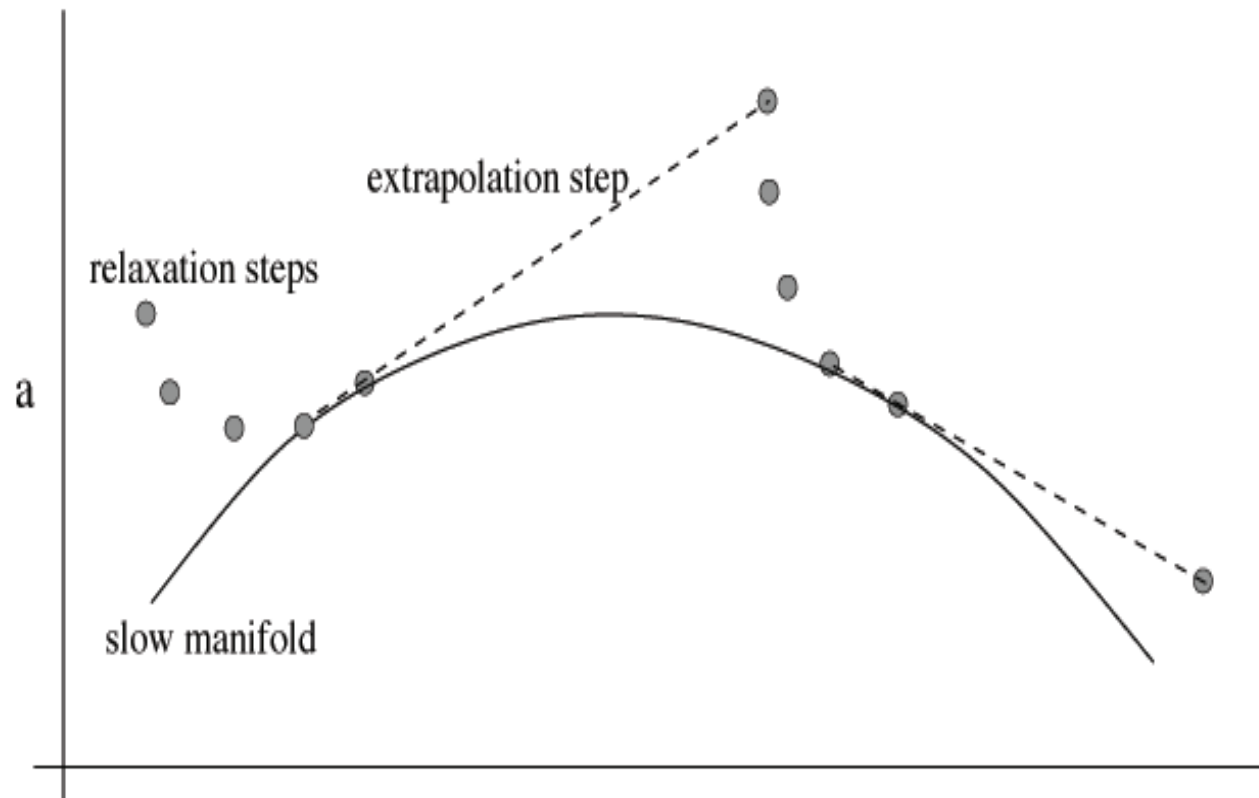


Figure 3: Schematic illustration of the projective integration scheme

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$$\begin{aligned}\frac{dx}{d\tau} &= \frac{1}{\epsilon} f_1(x, y) + f_0(x, x) \\ \frac{dy}{d\tau} &= \frac{1}{\epsilon^2} g_0(y, y) + \frac{1}{\epsilon} g_1(x, y)\end{aligned}$$

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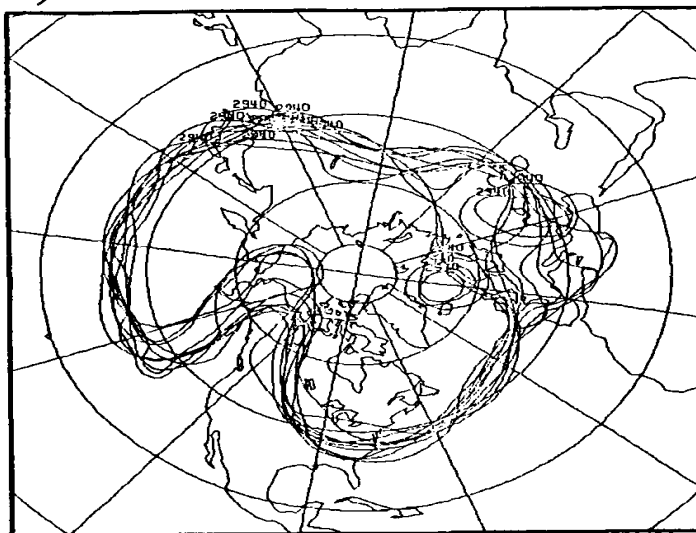
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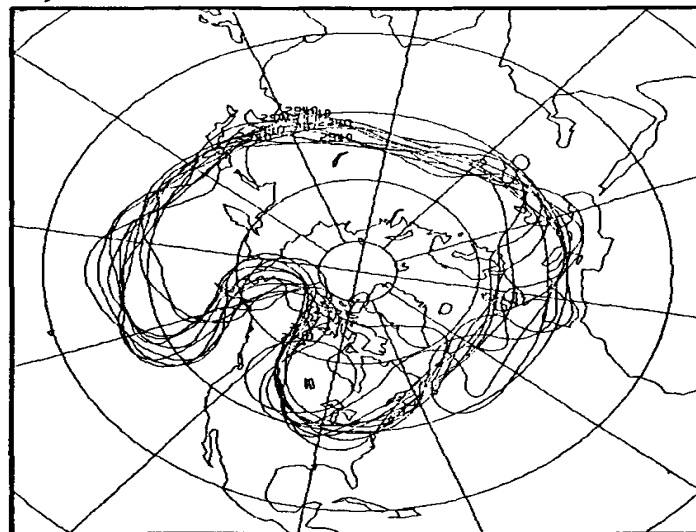
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 - ★ **the existence of limiting slow dynamics**

Atmospheric low-frequency variability (LFV)

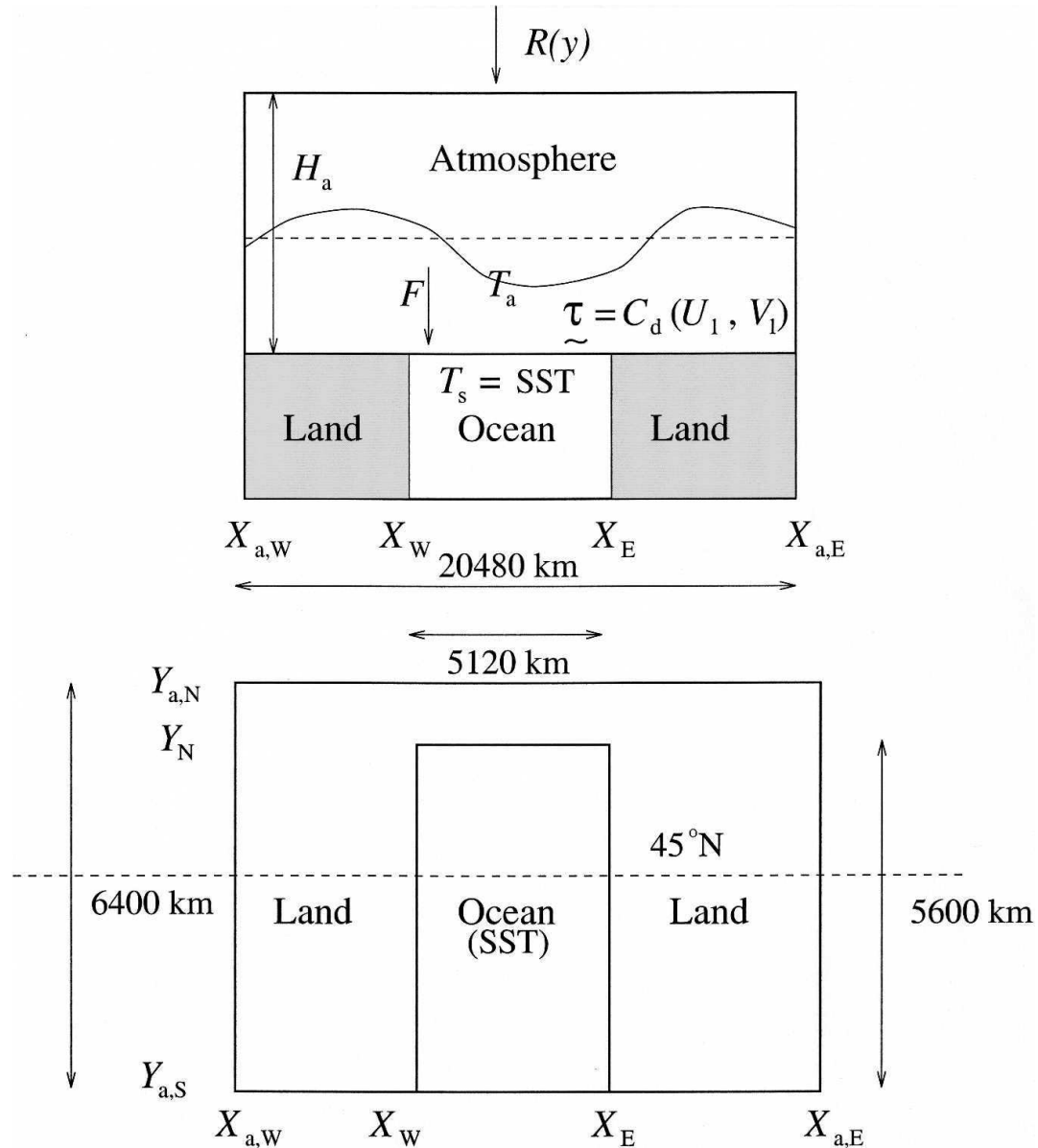
a) 2940M 61 1 20 - 61 1 29 0Z



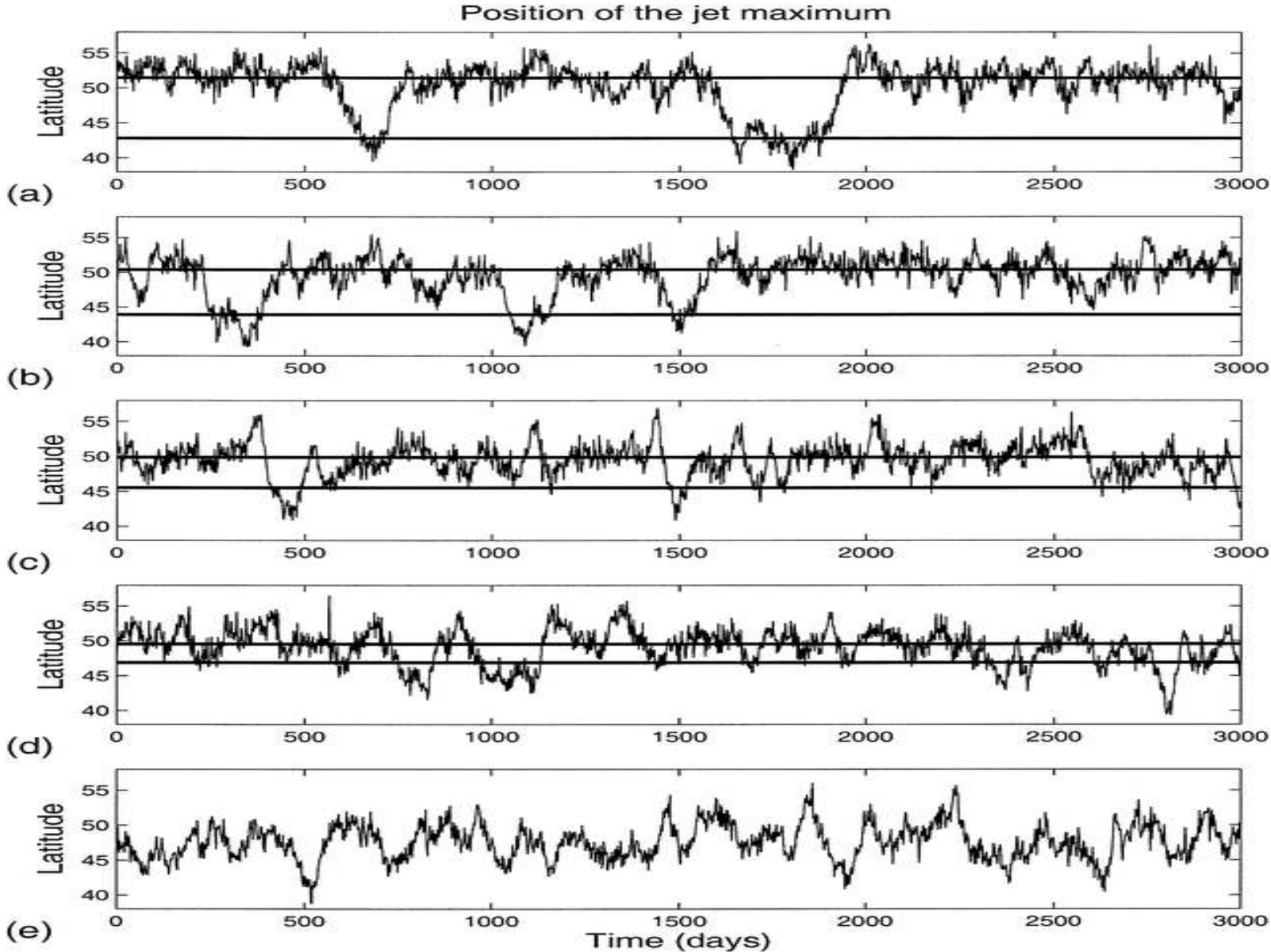
b) 2940M 78 12 31 - 79 1 9 0Z



QG model of LFV (Kravtsov et al. (2005))



Bifurcation: unimodal to bimodal distribution of jet axis



Existence of limiting slow dynamics

$$\begin{aligned}\frac{dx}{dt} &= f(x, y) \\ \frac{dy}{dt} &= \frac{1}{\delta\epsilon} g(x, y)\end{aligned}$$

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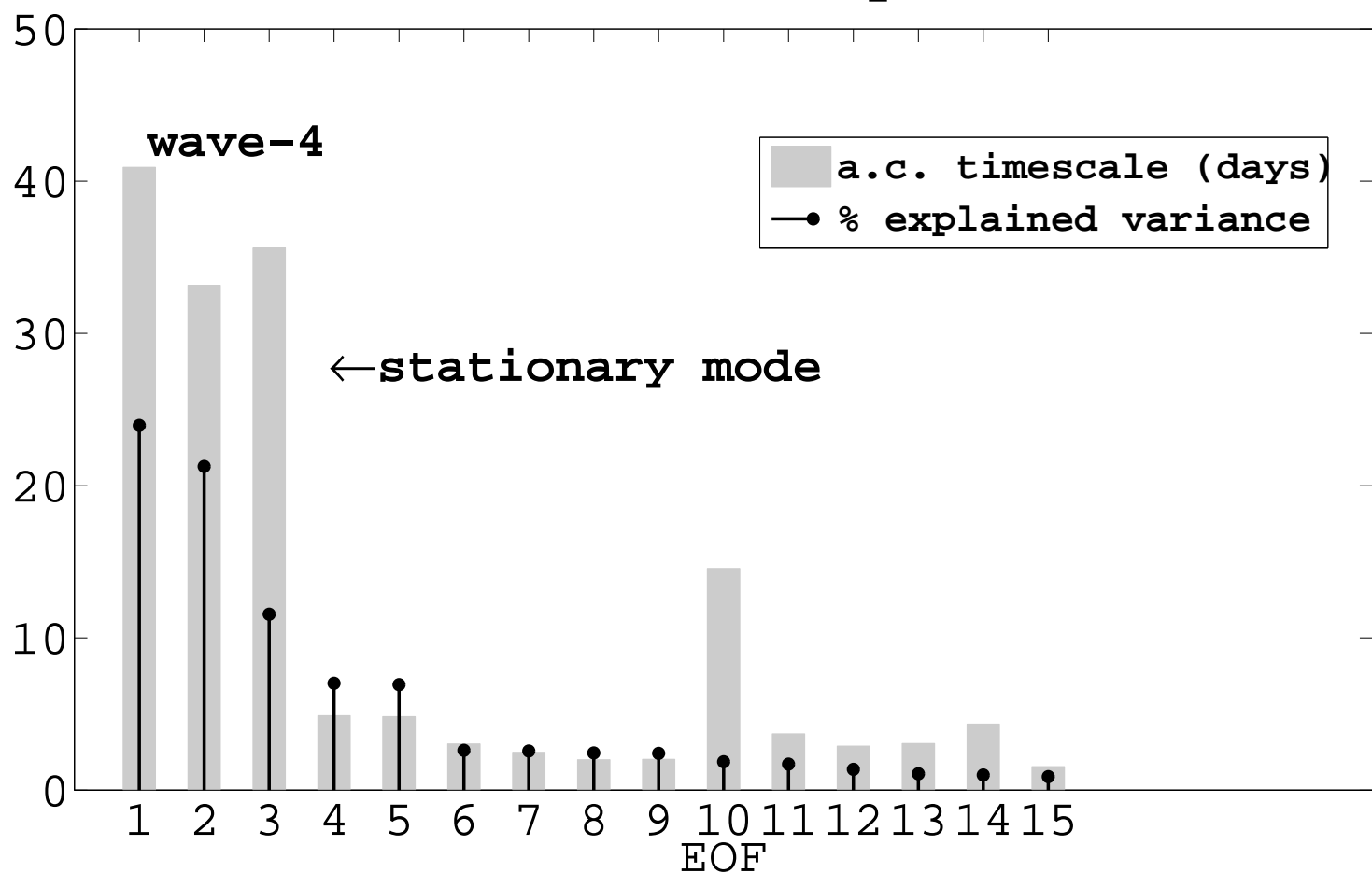
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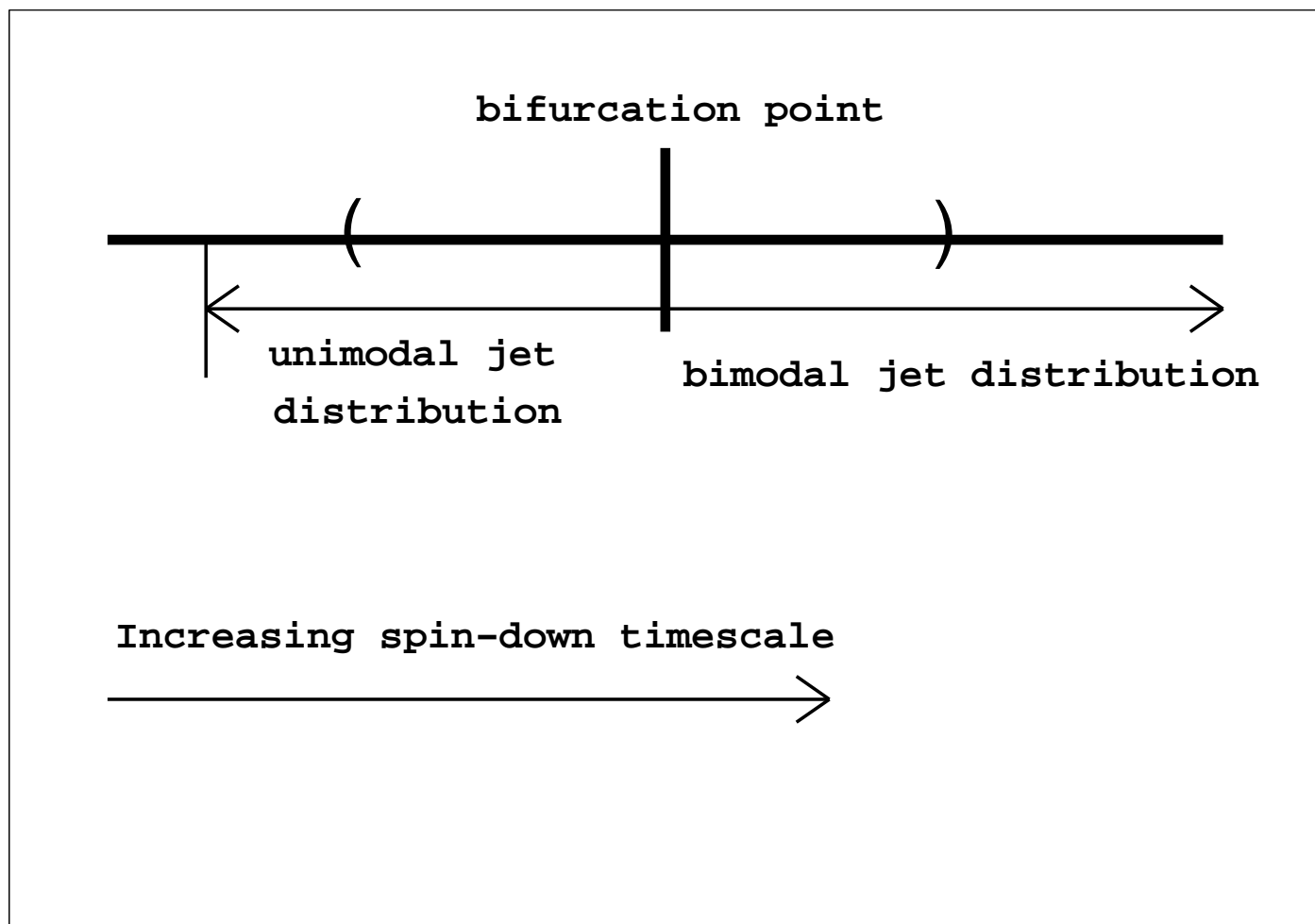
- The fast term is up-scaled by suitable choice of averaging parameters in Hasselmann's deterministic averaging equation
- Speeding up the fast mode is equivalent to changing the bifurcation parameter (the bottom drag parameter)
- In large neighbourhood of bifurcation point, **not every set of slow variables has limiting slow dynamics**

EOFs in region of jet bimodality

Autocorrelation timescale
and % explained variance vs. EOF;
 $k^{-1}=6.7$ days

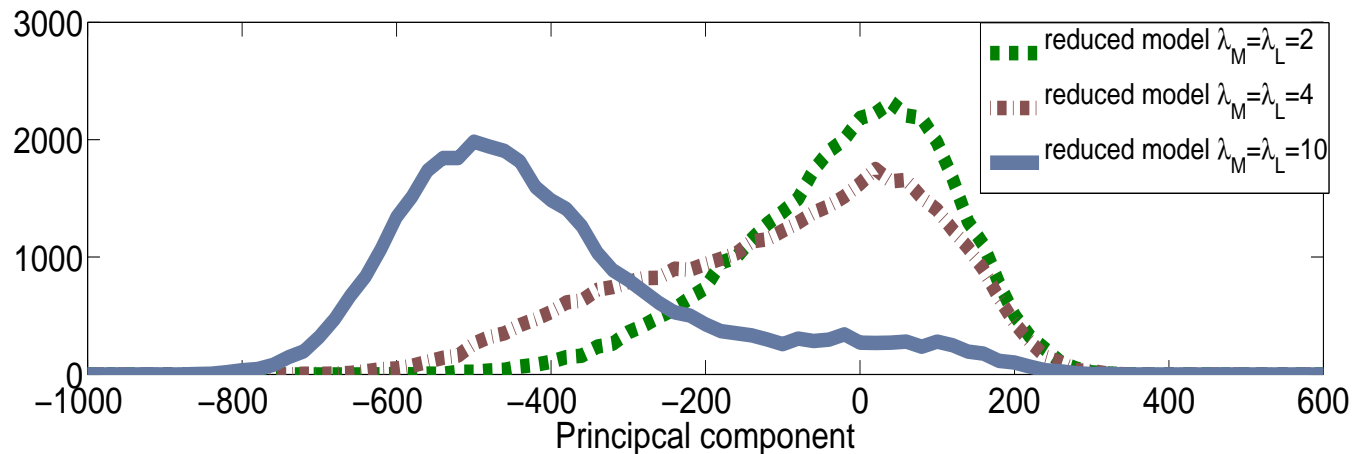
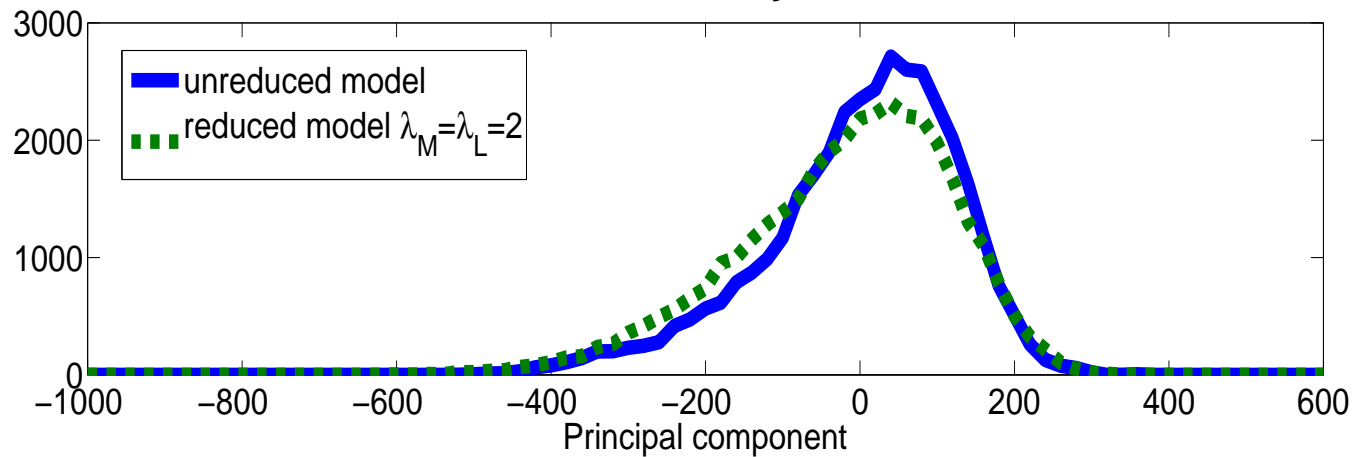


Schematic of KRG05 dynamics



Method of Franzke et al.(2005); 1-D SDE

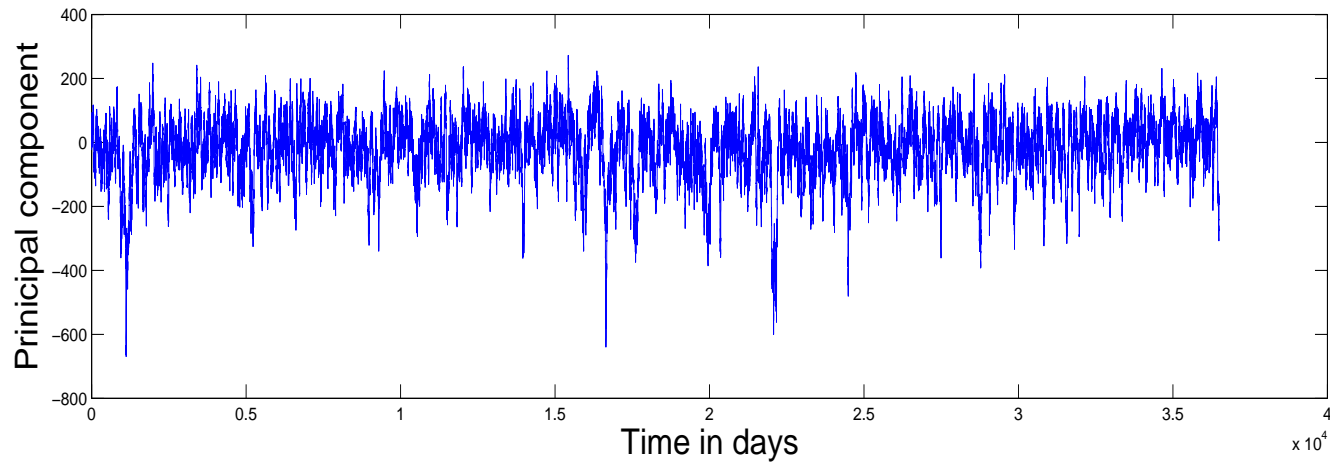
Stationary mode PDFs of unreduced and 1-D regressed reduced models;
 $k^{-1}=6.7$ days



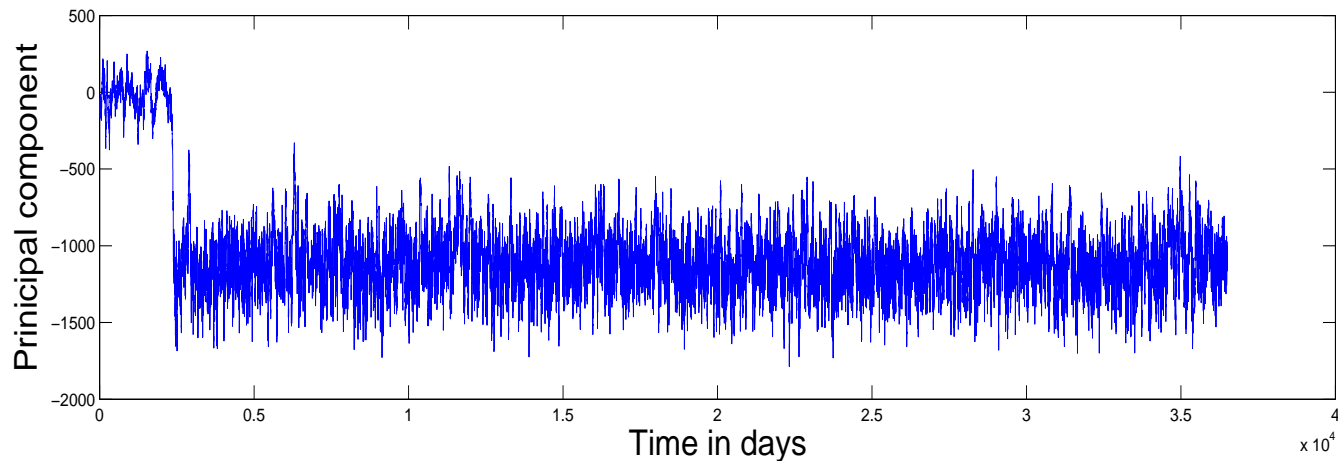
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Stationary mode time series of regressed reduced models w/out wave-4; $k^{-1}=6.7$ days

$$\lambda_B = \lambda_A = \lambda_F = 0.175$$

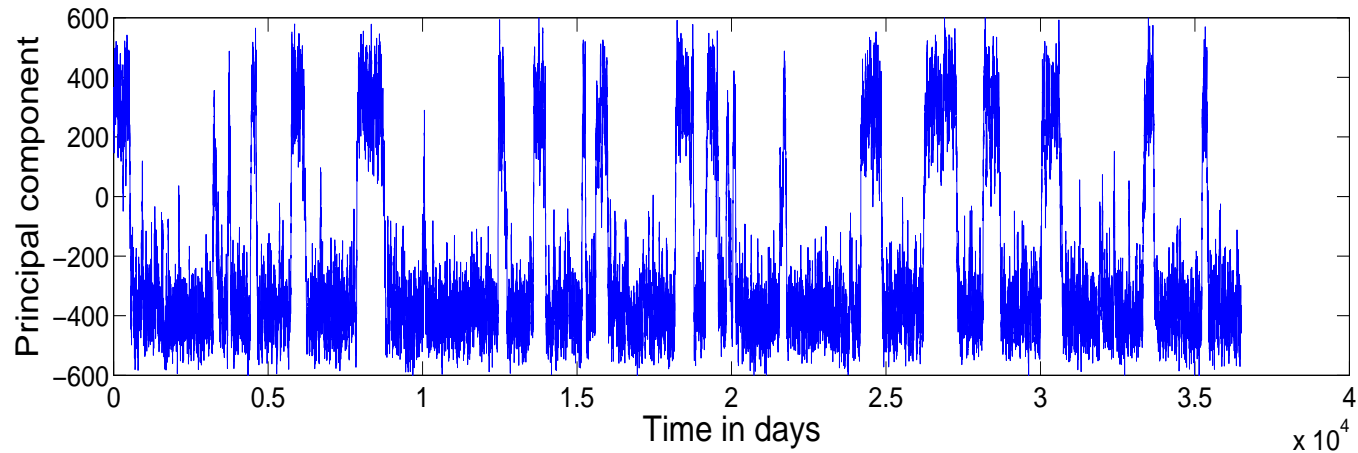


$$\lambda_B = \lambda_A = \lambda_F = 0.1$$

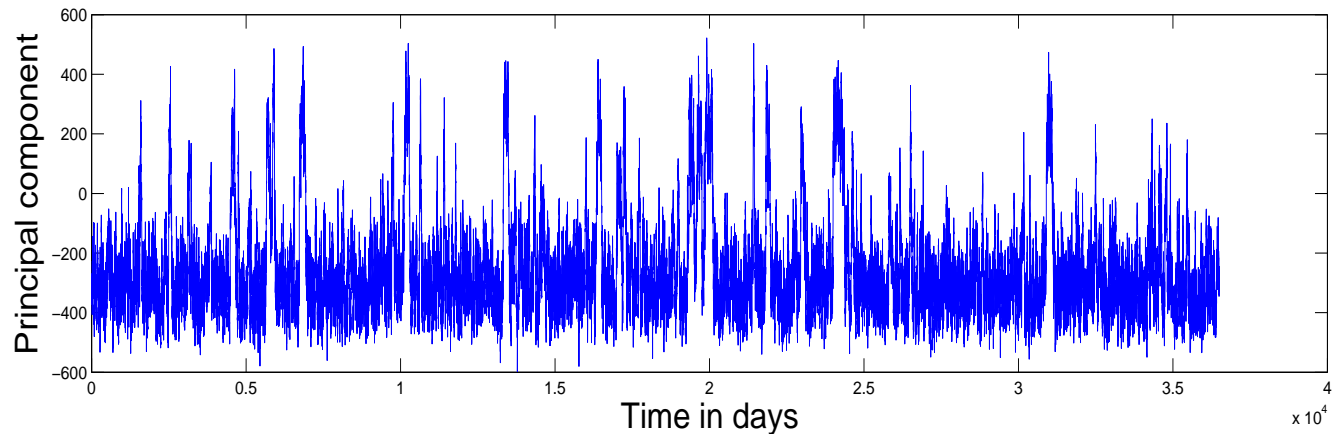


Method of Franzke et al.(2005); 3-D SDE

Time series of stationary mode; $k^{-1}=2.3$ days
Regressed reduced model with $\lambda_B=\lambda_A=\lambda_F=0$

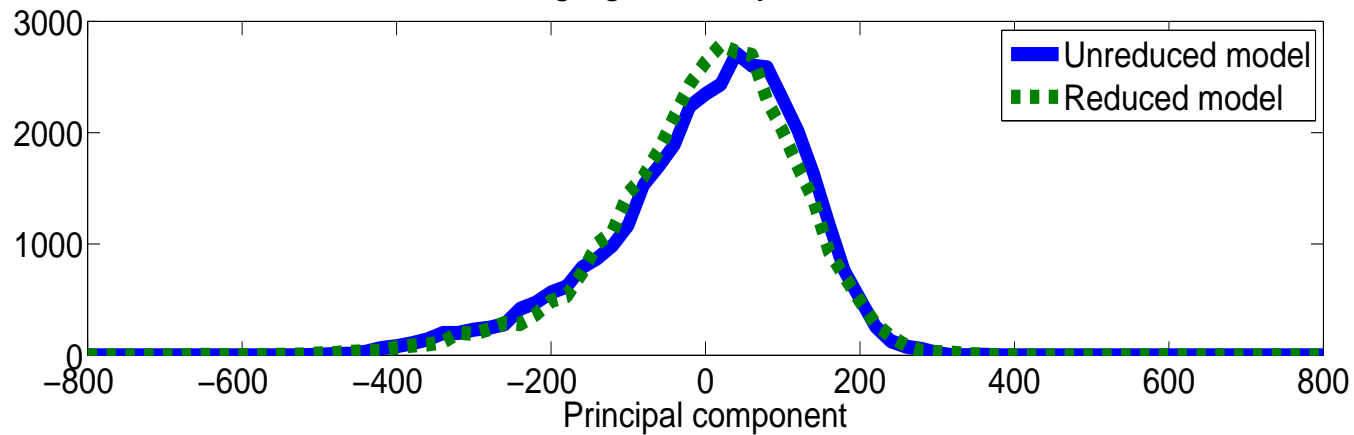


Regressed reduced model with $\lambda_B=\lambda_A=\lambda_F=0.15$

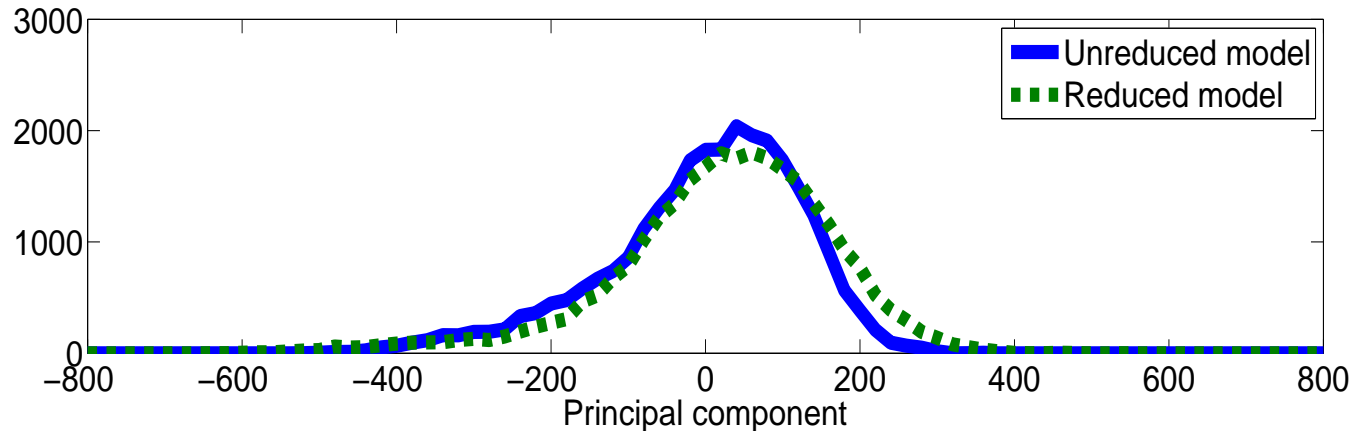


Hasselmann's method

Stationary mode PDFs of unreduced and reduced models; $k^{-1}=6.7$ days
Deterministic averaging–DNS hybrid; 4X faster than DNS



Deterministic averaging with additive white noise



Conclusions

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- The conclusion of KRG05 that first-order dynamics of jet bimodality arises from interaction between stationary and wave-4 modes is incorrect
- leading fast synoptic eddies are of first-order importance and wave-4 facilitates transitions between states

Conclusions: Method of Franzke et al. (2005)

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- For larger models, off-line calculations would be impractical without further simplifications

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